Myosin filament sliding through the Z-disc relates striated muscle fibre structure to function

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Striated muscle contraction requires intricate interactions of microstructures. The classic textbook assumption that myosin filaments are compressed at the meshed Z-disc during striated muscle fibre contraction conflicts with experimental evidence. For example, myosin filaments are too stiff to be compressed sufficiently by the muscular force, and, unlike compressed springs, the muscle fibres do not restore their resting length after contractions to short lengths. Further, the dependence of a fibre’s maximum contraction velocity on sarcomere length is unexplained to date. In this paper, we present a structurally consistent model of sarcomere contraction that reconciles these findings with the well-accepted sliding filament and crossbridge theories. The few required model parameters are taken from the literature or obtained from reasoning based on structural arguments. In our model, the transition from hexagonal to tetragonal actin filament arrangement near the Z-disc together with a thoughtful titin arrangement enables myosin filament sliding through the Z-disc. This sliding leads to swivelled crossbridges in the adjacent half-sarcomere that dampen contraction. With no fitting of parameters required, the model predicts straightforwardly the fibre’s entire force–length behaviour and the dependence of the maximum contraction velocity on sarcomere length. Our model enables a structurally and functionally consistent view of the contractile machinery of the striated fibre with possible implications for muscle diseases and evolution.

1. Background

As motors of life, muscles convert chemical energy into mechanical energy and heat. Muscles transport substances within the body, stabilize the skeleton and enable locomotion. The first property characterizing the mechanical function of striated muscles to be described [1] was the active isometric force–length relationship (FLR). This property shows the maximal forces the striated muscle fibre can produce by electrical stimulation at different constant lengths. The classic FLR (figure 1, straight lines) [2] with its strikingly linear segments has been described not only up to the fibre level but also for the whole muscle [7,8]. Despite the fact that this relationship represents basic textbook knowledge for life science students, to date a convincing structural model explaining the shape of the entire FLR does not exist.

According to the sliding filament [4,5] and crossbridge [6] theories, actin and myosin filaments slide relative to each other in response to forces generated by temporary crossbridges formed by myosin heads—projecting from the myosin filaments—and actin filaments. This yields straightforward geometric explanations for the plateau region and for the region of decreasing isometric force based on filament lengths (figure 1, straight lines) [2]. The slope change (figure 1, black circle) in the range of increasing isometric force (ascending limb) is typically related to myosin filaments hitting the Z-disc, a thin-meshed filament structure [9] defining the sarcomere boundary. It is, however, not clear which mechanism(s) are responsible for the decrease in force in the shallow and the steep slope regions of the ascending limb of the FLR [2,10–13].
More specifically, until now the ascending limb’s steep slope is typically explained with myosin filament folding or compression (figure 1, top left scheme), and a possibly decreasing number of crossbridges [2,10,13]. An alternative recent explanation considers two-dimensional force production of crossbridges [14]. Ignoring that myosin filaments reach the Z-disc at the point of slope change on the ascending limb (figure 2, circle), their model predicts that the increasing filament lattice spacing leads to a considerable reduction in longitudinal force, explaining the slope of the steep part of the ascending limb. However, several studies suggest a minor impact of lattice spacing compatible with the whole working range of the fibre on longitudinal force [15,16].

In conflict with the ideas of an internal counteracting compression spring or lattice spacing being responsible for the steep slope of the ascending limb, Ramsey & Street [3] reported a pronounced FLR foot region for short fibre lengths (figure 1, symbols) produced in experiments with prolonged stimulus duration. Several phenomena were observed during and after fibre contractions within this range. Tension developed very slowly at short fibre lengths and was lower after re-elongation of the fibres. Moreover, unlike their regular behaviour, fibres did not restore their resting length after cessation of activation. These observations cannot be explained by myosin filament compression, lattice spacing or by a decreasing number of available crossbridges, and hence seemingly contradict the classic sliding filament and crossbridge theories of muscle contraction.

Here we develop a model consolidating structural and micromechanical data that can explain the entire FLR and resolve the apparent conflicts between experimental evidence and classic theories of contraction. The few required model parameters are taken from the literature or set by structural arguments. The model suggests mechanical causes for the tetragonal arrangement [9,17] of actin filaments at the Z-disc (in contrast with their hexagonal arrangement in the typical actin–myosin filament overlap), for the so far unknown function of the second myosin head [18], and for the ability of myosin heads to form crossbridges with actin filaments of opposite polarity [19,20].

**2. Myosin filament sliding through the Z-disc**

The core of the model developed here is the idea that myosin filaments slide through the meshed Z-disc (figure 2 [21]) instead of being compressed or folded at short sarcomere lengths as typically assumed. This leads to several new types of actin–myosin filament overlap (figure 3b–c) affecting force production. In the following, we first introduce structural and geometrical arguments supporting the suggested mechanism and subsequently explain our model and its predictions (see Appendix A for mathematical details of the model).

The sliding of myosin filaments through the Z-disc requires a highly organized and sensibly adjusted geometry. First, the actin filament arrangement near the meshed Z-disc should be able to accommodate twice the number of myosin filaments within one half-sarcomere. The cross section of the typical actin–myosin overlap region reveals a hexagonal arrangement of actin filaments with a myosin filament centred in each actin hexagon (figure 3a, cross section *) [10,22]. Rearranging the actin filaments to a new regular lattice to accommodate twice the number of myosin filaments would require a tetragonal arrangement of actin filaments. Indeed, actin filament arrangement is tetragonal (square-cut) on either side of the meshed Z-disc (figure 3a, cross section *) [9,17].

Moreover, myosin filaments should be guided to form a regular pattern embedded in the tetragonal actin grid before reaching the Z-disc to enable orderly sliding through the Z-disc. Titin may be involved in fulfilling this task. Titin anchors myosin filaments to the Z-disc, and two titin
molecules connect to each actin filament within the Z-disc [24]. Because the ratio of actin to myosin filaments is 2 : 1 in regular overlap (figure 3a, cross section **), the ratio of titin molecules to myosin filaments is 4 : 1. A simple possible symmetry (accounting for the two titin molecules per actin filament) in the myofilament arrangement is when each myosin filament connects via four titin proteins to four actin filaments forming the vertices of a square, resulting in a centred chequered arrangement of myosin filaments within the tetragonal actin grid. This arrangement would facilitate highly organized myosin filament sliding through the Z-disc, with the myosin filaments possibly entering the empty squares within the actin filament grid of the adjacent half-sarcomere (figure 2).

Although a comparably dense structure [24], the Z-disc adapts in diameter when the muscle changes in length (e.g. [23]) and is essentially porous [17]. Actin filament ends of opposite polarity are cross-linked by α-actinin (35 nm long molecules), forming a flexible basket weave pattern in activated muscle [17,26]. This structure offers a sufficient number of channels required for myosin filaments passing the Z-disc (figure 2; electronic supplementary material, figure S1). However, preceding the proposed sliding, actin filament grids of opposite polarity from neighbouring sarcomeres are not aligned at the Z-disc but are shifted by half a grid cell (approx. 20 nm; electronic supplementary material, figure S1). Subtleties within the Z-disc structure [17] break the symmetry of the basket weave pattern, which may help to align the neighbouring actin filament grids during the sliding process (see electronic supplementary material, figure S1 and text S1). This may result in a myofilament configuration after sliding through the Z-disc similar to that shown in figure 2.

Tapered myosin filament ends [27] facilitate entering the meshed Z-disc and help to take advantage of subtle asymmetries mentioned above. Moreover, with shortening, the width of a sarcomere increases due to the preservation of volume in muscle [28]. Similarly, the Z-disc may expand laterally by unfolding its spatial zigzag structure facilitated by lower forces pulling on the actin filaments (figure 1), making room for the increasingly thicker myosin filaments sliding through the Z-disc. Despite the increased number of filaments in the developing overlaps (figure 3b–e), the lateral distance between actin and myosin filaments would be at least 90% of their distance at optimal length in the regular overlap arrangement (see electronic supplementary material, text S2).

3. The model and its predictions

In our half-sarcomere model, active isometric force is proportional to the effective actin–myosin overlap length where force can be generated by crossbridges. This effective length refers to the standard hexagonal arrangement of actin filaments surrounding each myosin filament (figure 3a, cross section **). Different overlaps (indicated by arrows in figure 3) contribute with different strength to force production. In the actin–actin–myosin overlap (figure 3a; the number of repetitions of the filaments’ names in the description of overlaps refers to the number of filaments in relation to regular overlap, i.e. actin–actin–myosin overlap contains twice the number of actin filaments and the identical number of myosin filaments like the regular overlap region), actin is abundant and seems to impede force production [13]. This result is extrapolated to the newly developing overlap in figure 3d, where actin filaments are also abundant. Because actin polarity determines the direction of action of
the crossbridges [19,20], half of the crossbridges in the remaining overlaps (figure 3b,c,e) would tend to shorten the sarcomere while the other half, the swivelled crossbridges [20], would tend to lengthen the sarcomere (figure 4). Assuming that the swivelled crossbridges produce the same force as their regular counterparts, the resulting isometric force in these overlaps equals zero, and the model generates an FLR (figure 5, prediction I) remarkably similar to the classic FLR (grey dotted lines).

However, the crossbridge is much softer in pushing than in pulling [29] (figure 4). From experimental data [29,30,33], we estimate (see electronic supplementary material, text S3) that the pushing force of a swivelled crossbridge in an isometric contraction is approximately half the pulling force of a regular crossbridge. Considering further the number of myosin filaments interacting with each actin filament for each overlap, the model predicts an FLR (figure 5, prediction II) that corresponds well with the classic experimental FLR, and—in the foot region below the length of classic zero force—with the data of Ramsey & Street [3] (figure 5, symbols).

4. Discussion

Our simple muscle model effortlessly explains the stunning linearity of the segments of the classic force–length relationship (FLR) of striated muscle (figure 5, prediction I). It is based on the idea that myosin filaments slide through the Z-disc [21] (figure 2) instead of being compressed or folded at the Z-disc as commonly assumed. By this elegant mechanism, the muscle can maintain its highly organized structure at short lengths (figure 5, range B) occurring regularly during daily activities [34]. Considering decreased swivelled relative to regular crossbridge forces, the model predicts as yet puzzling force production at lengths shorter than the length of classic zero force [3] (figure 5, prediction II, ranges C–E).

It seems noteworthy that the good agreement of the model predictions and data were achieved without fitting of parameters; instead, parameters were obtained from the literature or set by structural arguments (cf. Appendix A; electronic supplementary material, text S3). Decreased relative swivelled crossbridge force decreases the FLR’s ascending limb’s slope in our model (figure 5). This effect may be compensated by increasing filament lattice spacing, tending to increase the slope in this range [14]. Although direct micrographic evidence is scarce [21], a range of micro-mechanical, biophysical and structural evidence supports the theory.

(a) Flexural stiffness of myofilaments

The suggested mechanism relies on sufficient flexural stiffness of the myofilaments. The myosin filaments should not buckle when passing the Z-disc (figure 2) or when transmitting compressive force due to eccentric crossbridges (figure 3b–e). Conservative estimates of critical buckling force of myosin filaments based on experimental data (see electronic supplementary material, text S4) yield values of 0.90–1.33 nN, about three to four times the maximum isometric force in a myosin filament [35]. Moreover, muscle force decreases with contraction speed [36]. Hence myosin filament buckling...
seems rather unlikely; vice versa, myosin filament stiffness seems to be a crucial filament property, suggesting that myosin filaments cannot be folded or even compressed at the meshed Z-disc.

In comparison, actin filaments are much more flexible (approx. 70-fold) than myosin filaments [37]. Still, actin filaments are sufficiently stiff to pass the sarcomere’s M-line, and—after unphysiological detachment from the Z-disc—they can be pushed like rods to the opposing Z-disc by crossbridges acting at their ends [13]. Their flexural stiffness may become relevant within the range C (figure 5) when swivelled crossbridges would compress actin filaments (figure 3c, green myosin head on blue actin filament).

(b) Prediction of maximum contraction velocity

Otherwise constant maximum contraction velocity of sarcomeres in regions of zero passive force is known to decrease linearly for half-sarcomere lengths shorter than length $l_4$ [38] (figure 6, symbols). Swivelled crossbridges (figure 4c) formed when myosin filaments slide through the Z-disc into the adjacent half-sarcomere during half-sarcomere shortening slow down the contraction. Accounting for the force–velocity relation of muscle [36,39], our model even quantitatively predicts the reported behaviour of maximum contraction velocity in frog muscle (figure 6, prediction III; see electronic supplementary material, text S5), which supports the concept of myosin filament sliding through the Z-disc.

(c) ’Strange’ behaviour of muscle fibres

In early kinetic experiments on muscle fibres, Ramsey & Street [3] reported puzzling behaviour of muscle fibres at very short fibre lengths (figure 5, ranges C–E) and postulated that fibres enter a so-called ‘delta state’. The described slow development of force when entering range C coincides with the development of the new myosin–myosin–actin–actin overlap (figure 3c). Considering the flexural stiffness of the myofilaments (see above), it is unclear how this presumed overlap should seamlessly develop from the two different adjacent overlaps with triangular myosin filament lattice on one side (actin–actin–myosin overlap; figure 3a, cross section ***) and tetragonal myosin filament lattice on the other side (myosin–myosin–actin overlap; figure 3b). Myofilaments might interlock due to increased friction as a result of disordered filament sliding. This could be a contributing factor to the explanation of why fibres did not restore their resting length after cessation of activation [3]. If such a fibre is re-elongated, locked half-sarcomeres may, one by one, pop to a long length where passive half-sarcomere structures would take up the stretching force. In a subsequent end-isometric contraction, the fibre might exert less force than expected (as observed [3]) because of overstretched half-sarcomeres and compensating motions of half-sarcomeres accompanying force equilibration. Accordingly, if the whole fibres would be passively overstretched after entering the ‘delta state’, all stuck half-sarcomeres would elongate and they would restore their resting length, as seen in experiments [3].

Myosin-binding protein C, a thick filament protein that presumably bridges actin and myosin filaments [40], may also contribute to the described phenomenon of locked half-sarcomeres. Intriguingly, in frog myosin filaments this large protein occurs in 43 nm intervals from approximately 250 to 500 nm (measured from the M-line) along the myosin filament [40]. Hence, according to our model, the first layer of myosin-binding protein C would reach the Z-disc in range C (figure 5). It is feasible that over time myosin-binding protein C gets pulled through the Z-disc in activated muscle and then prevents passive re-elongation of the muscle to resting length when activation ceases (e.g. due to friction).

Although performed in part at unphysiological lengths, the measurements by Ramsey & Street [3] potentially contain important clues to the functioning of the contractile machinery of the muscle fibre in the physiological range. Because some doubt remains with respect to reproducibility and hence relevance of the findings by Ramsey & Street [3] obtained with frog muscle fibres for the general functioning of vertebrate muscle fibres, we repeated parts of their experiments with fibres of the rat M. extensor digitorum longus (figure 7). We likewise found considerable force development of the fibres for long duration of activation at lengths shorter (0.53 $\mu m$ half-sarcomere length) than classical zero force (0.63 $\mu m$ half-sarcomere length; figure 1 [2]). Also, fibres did not restore their resting lengths and produced less force when re-elongated, but full force after being passively overstretched. Moreover, Schoenberg & Podolsky [42] also measured forces at half-sarcomere lengths of 0.5–0.6 $\mu m$. There is increasing evidence that these results are not artefacts but properties of vertebrate fibres; myosin filament sliding through the Z-disc is an appealing possibility that reconciles them with current theories of contraction.

(d) Swivelled crossbridges and evolution

More than half a billion years ago, striated muscles presumably appeared in our prebilaterian ancestors [43]. Prior to the existence of exo- and endoskeletal structures, muscles probably adhered to gelatinous material such as is present,

Figure 6. Comparison of model-predicted maximum half-sarcomere contraction velocity (black line) and data (symbols [38]). The model and the data show a similar decrease of maximum contraction velocity for half-sarcomere lengths smaller than $l_4$. For orientation, the force–length relationship is depicted below. The half-sarcomere contraction velocity is normalized to $v_{\text{max}}$, its value in the plateau range of the force–length relationship. The increase in maximum contraction velocity in the grey shaded area stems from passive forces [38] not considered in our model.
for instance, in sponge and jellyfish [43]. Therefore, the muscle had to be equipped with a safety mechanism preventing extreme shortening and muscle malfunction, as described above. The swivelled crossbridges damping muscle contraction at short lengths represent such a safety mechanism. On a related note, this phenomenon may explain the so far unknown function of the second parallel myosin head attached to each myosin molecule [18], which may generate the swivelled crossbridges. However, this function comes at a metabolic cost. In species with skeletons, evolution may have adjusted the region of operation of the muscle by tuning origin and insertion such that it avoids damping its own contraction. For example, floating species such as fish can function more in the plateau region of the FLR [34] than non-floating species. In floating species, postural stability is provided by the aquatic environment. By contrast, species acting with their musculoskeletal system against gravity seem to make use of this ‘old’ muscle function by selected muscles operating in the range of the steep part of the ascending limb of the FLR [34], potentially to increase stability and to facilitate docile behaviour towards perturbations.

5. Conclusion
The assumption of myosin filament folding or even compression (which would need to exceed 60% for extreme cases of muscle shortening; figure 1) commonly used in textbooks cannot explain (i) how force can be produced at lengths shorter than the length of classic zero force; (ii) why the muscle does not restore its resting length after contractions with extended activation duration; and (iii) how the stiff myosin filaments can be folded or compressed at a meshed Z-disc. Accepting the idea that myosin filaments slide through the Z-disc and incorporating this mechanism into a comparably simple model with comprehensible, few parameters reconcile these data with classic theories of contraction.

In addition, the model enables new perspectives on the relation between striated fibre structure and its mechanical function. The transformation from hexagonal to tetragonal actin filament lattice near the Z-disc enables the development of an orderly myofilament structure at short fibre lengths. The Z-disc becomes an integral part of the contractile mechanism. For example, eccentric exercise [44] or muscle diseases like Duchenne muscular dystrophy or Sjögren’s syndrome are accompanied by a loss of structural integrity of the Z-disc. This may result in disordered myosin sliding through the Z-disc, contributing to the observed muscle weakness. A more comprehensible understanding of muscle structure and function seems within reach.

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Appendix A. The model

Active isometric half-sarcomere force $F$ depends on half-sarcomere length $h_0$ and is proportional to an effective actin–myosin overlap region $l_{\text{eff}}$, where force can be generated by crossbridges. Normalizing this length to the maximum effective actin–myosin filament overlap length ($l_{\text{eff,max}}$ the zone of myosin heads in a half-myosin), we obtain the isometric force as percentage of maximum isometric half-sarcomere force $F_{\text{im}}$ at optimal fibre length:

$$F(h_0) = \frac{l_{\text{eff}}(h_0)}{l_{\text{eff,max}}} F_{\text{im}}. \quad \quad (A1)$$

The Z-disc is assumed to have zero depth. Then, starting from long lengths, half-sarcomere lengths $l_1$ (no actin–myosin overlap), $l_2$ (start of plateau), $l_3$ (start of shallow slope), $l_4$ (start of steep slope), $l_5$ (actin–actin–myosin overlap meets myosin–myosin–actin overlap), $l_6$ (actin filaments of adjacent half-sarcomeres meet Z-disc), $l_7$ (myosin filaments of adjacent half-sarcomeres meet bare zone) and $l_8$ (myosin filaments of adjacent half-sarcomeres meet M-line) are given by

$\begin{align*}
  l_1 &= l_a - \frac{l_{m}}{2}, \\
  l_2 &= l_a + \frac{h_{\text{bare}}}{2}, \\
  l_3 &= l_a - \frac{h_{\text{bare}}}{2}, \\
  l_4 &= \frac{l_m}{2}, \\
  l_5 &= \frac{2l_a + l_m}{6}, \\
  l_6 &= \frac{l_m}{4}, \\
  l_7 &= \frac{l_m + h_{\text{bare}}}{4}, \\
  l_8 &= \frac{l_m - h_{\text{bare}}}{4},
\end{align*}$

$$\quad \quad (A2)$$

with actin filament length $l_a$, myosin filament length $l_m$ and the bare zone length without myosin heads at the centre of the myosin filament $h_{\text{bare}}$. Each different overlap region (figure 3, between arrows) is scaled by an individual factor $c$ to a corresponding length of regular overlap (i.e. three myosin filaments interact with each actin filament tending to shorten the sarcomere); the sum of the scaled lengths yields the effective overlap length

$$l_{\text{eff}}(h_0) = \begin{cases} 
(l_1 - h_0), \\
(l_1 - l_2), \\
(l_m/2 + h_0 - l_a) + c_{aam}(l_3 - h_0), \\
3h_0 - l_a - l_m/2 + c_{aam}(l_3 - h_0) + c_{aamm}(l_m/2 - h_0), \\
c_{aamm}(2h_0 - l_m/2 - h_{\text{bare}}/2) + c_{aamm}(2h_0 - l_a) + c_{aam}(l_a + l_m/2 - 3h_0), \\
c_{aamm}(2h_0 - l_m/2 - h_{\text{bare}}/2) + c_{aamm}(l_m/2 - h_0 - l_a) + c_{aammm}(l_a - 2h_0), \\
c_{aammm}(3h_0 - l_a - l_m - h_{\text{bare}}/2) + c_{aammm}(l_a - 2h_0) + c_{aamm}(2l_7 - 2h_0).
\end{cases} \quad \quad (A3)$$

The index of $c$ denotes the corresponding overlap region. For instance, $c_{aamm}$ scales the actin–actin–myosin–myosin overlap (figure 3c). The asterisk in the bottom row of equation (A3) indicates its quantitative change within the bare zone of myosin of the half-sarcomere of interest where only half of myosin heads are available (figure 3e).

Calculating prediction I (figure 5), all $c$ factors equal zero. For calculating prediction II (figure 5) and prediction III (figure 6), $c$ factors consider the available number of actin filaments, the number of myosin filaments interacting with each actin filament and the orientation of the crossbridges. $c_{aamm}$ and $c_{aamm}$ were set to zero because experiments indicate that force is zero when actin filaments are abundant [13]. From the cross sections in figure 3b,c,e, we count how many myosin filaments produce regular ($M_f$) and swivelled crossbridges ($M_s$) with each actin, respectively. Sectors $M_s$ tending to lengthen the sarcomere produce half the force ($f = 0.5$; see electronic supplementary material, text S3) of filament sectors $M_f$ tending to shorten the sarcomere. Thus, we subtract the number of $M_s$ per actin filament scaled by $f$ from $M_f$ per filament and divide this difference by three (the number of myosin filaments per actin tending to shorten the sarcomere in regular overlap). Finally, this number is multiplied by $A$ (the number of available actin filaments divided by the number of actin filaments in regular overlap) to yield individual $c$:

$$c = A \cdot \frac{M_f - f \cdot M_s}{3}. \quad \quad (A4)$$

According to the mechanism of myosin sliding through the Z-disc (figure 2), we expect a regular tetragonal pattern of actin filaments and myosin filaments within the myosin–myosin–actin overlap (figures 2 and 3b, cross section). Within this overlap region, two myosin filaments from the current half-sarcomere and two myosin filaments from the adjacent half-sarcomere (forming swivelled crossbridges) act on each actin filament. Hence, $c_{aamm} = (1 - 0.5 - 2)/3 = 0.33$. In the actin–actin–myosin–myosin overlap the actin–myosin filament ratio is 2:1. Though formation of a regular structure may be inhibited (see Discussion), for the sake of
simplicity we assume the development of a hexagonal actin filament arrangement similar to the regular overlap zone. Here, on average, 1.5 myosin filaments per actin filament tend to shorten the sarcomere, and 1.5 myosin filaments per actin filament tend to extend the sarcomere. Because there are twice as many actin filaments compared with regular overlap, this results in $c_{aann} = 2 \times (1.5 - 0.5 \times 1.5)/3 = 0.5$. Accordingly, the actin–actin–myosin–myosin overlap within the bare zone of one set of myosin filaments must be scaled by $c_{aann}^{**} = 0.25$.

References


